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# Permittivity measurements at millimeter wave frequencies using dielectric rod resonator excited by NRD-guide

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#### Abstract

A method of measuring the relative complex permittivity ( $\varepsilon_r = \varepsilon' - j\varepsilon''$ , tan  $\delta = \varepsilon''/\varepsilon'$ ) for low-loss dielectric materials at millimeter wave frequencies has been developed, using a dielectric rod resonator excited by the nonradiative dielectric waveguide (NRD-guide). Relative permittivity ( $\varepsilon'$ ) and loss factor (tan  $\delta$ ) of the rod specimen are determined by the resonant frequency ( $f_0$ ) and unloaded Q-factor ( $Q_u$ ) of a TE<sub>0m1</sub> mode resonator. The effective conductivity ( $\sigma$ ) of conducting plates for short-circuiting the rod resonator is determined using TE<sub>021</sub> and TE<sub>02 $\delta$ </sub> mode sapphire resonators. Temperature dependence of  $\varepsilon'$  and tan  $\delta$  of sapphire and cordierite ceramics were evaluated at 60 GHz. This method has been adopted as the Japanese Industrial Standard (JIS R 1660-3) and is being prepared for the IEC international standard. Several standardized specifications are presented.

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### 1. Introduction

We need a reliable method for measuring  $\varepsilon_r$  of the dielectric materials for designing devices used in millimeter wave communication or sensing systems. Several methods<sup>1–4</sup> have been developed for measuring  $\varepsilon_r$  of the low-loss dielectric materials at the millimeter wave frequencies. Among these methods, a dielectric rod resonator method exited by the NRD-guide<sup>5</sup> has an advantage in that excitation of resonance is very easy. The reason is that a dominant LSM mode of the NRD-guide can be easily coupled to the TE mode of the dielectric resonator.

We have developed a simple and accurate method for measuring  $\varepsilon_r$  at millimeter wave frequencies based on the dielectric rod resonator excited by NRD-guide.<sup>6</sup> The values of  $\varepsilon'$  and tan  $\delta$ are determined using the TE<sub>0m1</sub> (m = 1, 2, 3) mode resonator. The TE<sub>0m1</sub> mode resonator allows us to achieve stable measurements, and accurate and simple calculations for  $\varepsilon'$  and tan  $\delta$  by analytic expression. The effective conductivity  $\sigma$  of the conducting plates is determined using TE<sub>021</sub> and TE<sub>02 $\delta$ </sub> mode sapphire resonators.

#### 2. Theory and measurement formulas

#### 2.1. Relative permittivity $\varepsilon'$ and loss factor tan $\delta$

The values of  $\varepsilon'$  and tan  $\delta$  of the rod specimen are determined using the TE<sub>0m1</sub> mode resonator, as shown in Fig. 1a for m = 2. A dielectric rod with diameter (*d*) and height (*h*) is short-circuited at both ends by two parallel conducting plates.

The values of  $\varepsilon'$  is calculated from the measured  $f_0$ , d and the space ( $h_c$ ) between the upper and lower conducting plates:

$$\varepsilon' = \left(\frac{\lambda_0}{\pi d}\right)^2 (u^2 + v^2) + 1 \tag{1}$$

where

$$v^{2} = \left(\frac{\pi d}{\lambda_{0}}\right)^{2} \left[ \left(\frac{\lambda_{0}}{2h_{c}}\right)^{2} - 1 \right]$$
(2)

Here,  $\lambda_0 = c/f_0$  is the free space resonance wavelength and *c* is the velocity of light. A formula for *u* is given elsewhere.<sup>6,7</sup>

Next,  $\tan \delta$  is calculated from the measured  $Q_u$ :

$$\tan \delta = \frac{A}{Q_{\rm u}} - BR_{\rm s} = \frac{A}{Q_{\rm u}} - \frac{B'}{\sqrt{\sigma_{\rm r}}}$$
(3)

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Fig. 1. Configuration of dielectric resonator for measuring  $\varepsilon_r$  and  $\sigma_r$ . Value of  $\varepsilon_r$  is measured using (a) and  $\sigma_r$  is measured using (a) and (b).

where

$$R_{\rm s}(\Omega) = \sqrt{\frac{\pi f_0 \mu}{\sigma}} = \sqrt{\frac{\pi f_0 \mu}{\sigma_0 \sigma_{\rm r}}} \tag{4}$$

Here,  $R_s$  and  $\sigma$  are the surface resistance and the effective conductivity of conducting plates, respectively. The relative conductivity is defined as  $\sigma_r = \sigma/\sigma_0$  ( $\sigma_0 = 5.8 \times 10^7$  S/m). Furthermore,  $\mu$  is the permeability and  $\mu = \mu_0 = 4\pi \times 10^{-7}$  for non-magnetic conducting plates. Formulas for *A* and *B* are given elsewhere.<sup>6,7</sup>

# 2.2. Determination of relative conductivity $\sigma_r$ of conducting plates

The value of  $\sigma_r$  of the conducting plates must be accurately measured, to determine tan  $\delta$  of low-loss dielectric specimen by (3). The TE<sub>021</sub> and TE<sub>02 $\delta$ </sub> mode dielectric rod resonators in Fig. 1, called "standard resonators", are used for measuring  $\sigma_r$ . Each standard resonator is made of sapphire single crystals with the same  $\varepsilon_r$ . They are designed to have the same  $f_0$ .

The standard TE<sub>02 $\delta$ </sub> resonator has a large dimension ratio (*d/h*). The conductor loss of the TE<sub>02 $\delta$ </sub> resonator is larger than that of the TE<sub>02 $\delta$ </sub> standard resonator, since the electromagnetic field in the TE<sub>02 $\delta$ </sub> resonator is concentrated near the surface of the lower conductor. High accuracy measurement of  $\sigma_r$  is achieved by enlarging the difference in the conductor loss of the two resonators. While  $f_{01} = f_{0\delta}$ ,  $Q_{u1}$  is higher than  $Q_{u\delta}$ , where subscripts 1 and  $\delta$  denote each standard resonator.

The value of  $\sigma_r$  can be calculated from the measured  $f_0$  (= $f_{01} = f_{0\delta}$ ),  $Q_{u1}$  and  $Q_{u\delta}$ :

$$\sigma_{\rm r} = \frac{\sigma}{\sigma_0} = \pi \mu f_0 \left[ \frac{Q_{\rm u1} Q_{\rm u\delta}}{G_1 G_\delta} \frac{G_1 P_{\rm e1} - G_\delta P_{\rm e\delta}}{Q_{\rm u1} P_{\rm e1} - Q_{\rm u\delta} P_{\rm e\delta}} \right]^2 / \sigma_0 \tag{5}$$

where partial electric energy filling factors  $(P_{e1})$  and  $(P_{e\delta})$ , and geometric factor  $(G_1)$  and  $(G_{\delta})$  of the resonators are defined elsewhere.<sup>6</sup> An example of their values is shown in Table 1.

Furthermore,  $\tan \delta$  of the standard resonators can be calculated:

$$\tan\delta_1 = \tan\delta_\delta = \frac{1}{Q_{u1}Q_{u\delta}} \times \frac{G_1Q_{u\delta} - G_\delta Q_{u1}}{G_1P_{e1} - G_\delta P_{e\delta}}$$
(6)

#### 3. Preparation of dielectric specimen

Typical specifications of the standard rod resonators are described. Each rod consists of the sapphire with low tan  $\delta$ . The axis of the each rod is parallel to the *C*-axis of sapphire. Table 1 shows *d*, *h*, *h*<sub>c</sub>, *P*<sub>e</sub> and *G* of the standard sapphire rods with  $\varepsilon' = 9.40$  perpendicular to the *C*-axis, for measuring  $\sigma_r$  at 60 GHz. The factors *P*<sub>e1</sub> and *G*<sub>1</sub> are related with *A* and *B* in (3) by *P*<sub>e1</sub> = 1/*A* and *G*<sub>1</sub> = *A*/*B*. In contrast, calculations for factors *P*<sub>e\delta</sub> and *G*<sub>\delta</sub> obtained by axis symmetric FEM calculations.

The TE<sub>011</sub>, TE<sub>021</sub> and TE<sub>031</sub> mode rod resonators are used to measure  $\varepsilon'$  and tan  $\delta$ , for materials with  $\varepsilon' = 2-4$ ,  $\varepsilon' = 4-20$ and  $\varepsilon' = 20-30$ . Fig. 2 shows desirable values of diameter (*d*) of the rod specimen for 60 GHz measurement as a function of  $\varepsilon'$ .

#### 4. Measurement apparatus

Two types of apparatus were used, as shown in Fig. 3. The dielectric rod resonator was coupled equally to the input and output NRD-guide. Space between the rod and the NRD-guide was adjusted so that insertion loss ( $IL_0$ ) at  $f_0$  was 20–25 dB. Dielectric strips of the NRD-guide with a width of 2.00 mm were made of Rexolite-1422 (cross-linked styrene copolymer) with  $\varepsilon' = 2.5$ ,

Table 1

Dimensions, partial electric energy filling factor  $P_{\rm e}$  and geometric factor G for standard sapphire rods for measuring  $\sigma_{\rm r}$  at 60 GHz.

Space, $h_c$ (mm)	TE <sub>021</sub>				TE <sub>02δ</sub>			
	Diameter, d (mm)	Height, <i>h</i> (mm)	Pel	$G_{1}\left( \Omega  ight)$	d (mm)	h (mm)	$P_{e\delta}$	$G_{\delta}\left( \Omega ight)$
2.25	3.14	2.20–2.25	0.915	1182	4.49	0.80	0.906	409



Fig. 2. Diameter *d* of dielectric specimen for measuring  $\varepsilon_r$  at 60 GHz, for height h = 2.20-2.25 mm and space  $h_c = 2.25$  mm.

for 60 GHz measurements. The apparatus had transducers from the NRD-guide to the waveguide. The end of the dielectric strip was sharpened in the transducer.

A-type apparatus was used for measuring the temperature dependence of tan  $\delta$ , since the TE<sub>02 $\delta$ </sub> standard resonator is easily constituted. The upper conducting plate was inset in the con-

ductors of NRD-guide in the A-type apparatus. The apparatus had a small air gap (g) between the upper conducting plate and the dielectric specimen. The calculations of  $\varepsilon'$  by (1) and of tan  $\delta$  by (3) are accurate for  $g < 50 \,\mu\text{m}$ . The error of  $\Delta \varepsilon'/\varepsilon'$  for  $g = 50 \,\mu\text{m}$  is about 0.01% and is negligibe.<sup>6</sup> In contrast, B-type apparatus was used for measuring the temperature dependence of  $\varepsilon'$ . The upper conducting plate was put on the specimen in this apparatus.

The apparatus was connected to a scalar network analyzer HP-8757 system. The  $Q_u$  was calculated from  $f_0$ , the half-power band width  $f_H - f_L$  and  $IL_0$ :

$$Q_{\rm u} = \frac{f_0 / (f_H - f_L)}{(1 - 10^{-IL_0/20})} \tag{7}$$

# 5. Results

The temperature (*T*) dependence of  $\varepsilon_r$  of sapphire and cordierite ceramics<sup>8</sup> was evaluated at 60 GHz by this method. First,  $f_0$  was measured as a function of *T* using the B-type apparatus. Then  $\varepsilon'$  was calculated from the  $f_0$ . Fig. 4a shows  $\varepsilon'$  and  $f_0$  of the TE<sub>021</sub> standard sapphire rod ( $d=3.130\pm0.005$  mm,  $h=2.250\pm0.001$ ,  $\alpha=5.8$  ppm/°C). Here,  $\alpha$  is the coefficient of thermal expansion. The values  $\varepsilon'$  of the sapphire increased linearly with increasing *T*. The first and second





Fig. 3. Measurement apparatus.



Fig. 4. Temperature dependence measurements of  $\varepsilon_r$  of sapphire crystals and cordierite ceramics, and of  $\sigma_r$  of conducting plates.

measurements were in agreement. The temperature coefficient of permittivity ( $TC\varepsilon$ ) of the sapphire was calculated to be  $86.4 \pm 0.8 \text{ ppm/}^{\circ}\text{C}$ , which was in good agreement with a result reported elsewhere.<sup>9</sup> Fig. 4b shows measurements of cordierite ceramics rod (d=4.803, h=2.253,  $\alpha=0.5$ ). The values  $\varepsilon'$  of the cordierite rod also increased linearly with increasing *T*. The value  $TC\varepsilon$  was calculated to be  $56.9 \pm 0.3 \text{ ppm/}^{\circ}\text{C}$ .

Next, to determine  $\sigma_r$  of the conducting plates of Cu and tan  $\delta$  of the standard sapphire,  $Q_{u1}$  and  $Q_{u\delta}$  of standard sapphire resonators was measured twice against *T* as shown in Fig. 4c, using the A-type apparatus with  $H_c = 2.279$  mm. Fig. 4d and 4e show  $\sigma_r$  of the conducting plates and tan  $\delta$  of the standard sapphire calculated by  $Q_{u1}$  and  $Q_{u\delta}$ . The value of  $\sigma_r$  decreased with increasing *T*. The value  $\sigma_r$  was 87.4% at 20 °C. The value tan  $\delta$  of the standard sapphire increased with increasing *T*. The value  $f_0/\tan \delta = 1.04 \times 10^6$  GHz at 20 °C was in agreement with a result reported elsewhere.<sup>6</sup> Furthermore, Fig. 4f shows  $Q_u$  and tan  $\delta$  of cordierite ceramics. This variation of tan  $\delta$  with *T* was relatively small.

# 6. Conclusion

A method of measuring  $\varepsilon_r$  at millimeter wave frequencies has been developed, using a dielectric resonator excited by the NRD-guide. Typical specifications of the rod specimen for  $\varepsilon_r$ measurements and of the standard sapphire rods for measuring  $\sigma_r$  at 60 GHz were presented. The temperature dependence of  $\varepsilon_r$  of sapphire and cordierite ceramics was accurately evaluated at 60 GHz by this method. The repeated measurements showed the error of  $TC\varepsilon$  was less than 1 ppm/°C.

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